## PHYS 301 HOMEWORK #9-- Solutions

1. For this problem, it is important to remember that if a vector field is conservative (i.e., its curl = 0), then the vector can be derived from a scalar potential. Additionally, if the field is conservative, its line integral is path independent, and can be computed either directly as a line integral, or by finding the value of its potential at the end points. For the two vectors given, a) is non - conservative (it should be easy to show its curl is non - zero) and b) is conservative.

Therefore, we have to compute the line integral for case a) explicitly for the two paths given. Recall that :

$$\int_{C} \mathbf{F} \cdot d\mathbf{l} = \int_{C} (F_x \, dx + F_y \, dy + F_z \, dz)$$

For the vector in case a) :

$$F_x = x y \qquad F_y = 2 y z \qquad F_z = 3 x z$$

so the line integral becomes :

$$\int_{C} (x y dx + (2 y z dy + 3 x z dz))$$

The first path is the straight line y = x, so we can parameterize this integral as x = y = t; dx = dy = dt. Since we are in the xy plane, z = dz = 0, and the line integral is simply :

$$\int_0^1 t^2 \, dt = \frac{1}{3}$$

Along the path  $y = x^2$ , we set:

$$x = t$$
;  $dx = dt$ ;  $y = t^2$ ;  $dy = 2t dt' z = dz = 0$ 

and the line integral becomes :

$$\int_0^1 t^3 \, dt = \frac{1}{4}$$

The second vector field is conservative, so we should expect the line integral to be the same for both paths (but let's do both to verify). Using the parameterization x = y = t; dx = dy = dt; z = dz = 0:

$$\int_0^1 (t^2 \, dt \, + \, 2 \, t^2 \, dt) \, = \, 1$$

Along the curved path, we set :

$$x = t dx = dt y = t^2 dy = 2t dt z = dz = 0$$

and the line integral becomes :

$$\int_0^1 (t^4 dt + 2 t^3 (2 t dt)) = \int_0^1 5 t^4 dt = 1$$

as we expect.

We can also find the potential from which this force is derived and find the work by evaluating the potential at the end points of the path. We can write :

$$\mathbf{F}_{2} = \nabla \phi \Rightarrow \left\{ y^{2}, 2 x y + z^{2}, 2 y z \right\} = \left\{ \frac{\partial \phi}{\partial x}, \frac{\partial \phi}{\partial y}, \frac{\partial \phi}{\partial z} \right\}$$

Equating the x components of each vector and integrating gives us :

$$\frac{\partial \phi}{\partial x} = y^2 \Rightarrow \phi = y^2 x + g(y, z)$$

where the constant of integration is constant with respect to x, so to be as general as we can be we write the constant as a function of y and z. If we differentiate this expression for  $\phi$  with respect to y, and equate it to the y component of the vector we obtain :

$$\frac{\partial \phi}{\partial y} = 2yx + g'(y, z) = 2xy + z^2$$

This means that g' (y, z) =  $z^2$ . Integrating this tells us g(y,z) =  $z^2$ y + h (z). Substituting this back into our expression for  $\phi$  gives:

$$\phi = x y^2 + y z^2 + h (z)$$

Now, we differentiate  $\phi$  with respect to z and equate with the z component of the vector :

$$\frac{\partial \phi}{\partial z} = 2yz + h'(z) = 2yz \Rightarrow h'(z) = 0 \text{ and } h(z) = \text{numerical constant}$$

Thus, apart from a numerical constant, we know that

$$\phi = x y^2 + y z^2$$

Finally, we compute the work from :

$$W = \phi(1, 1, 0) - \phi(0, 0, 0) = 1$$

2. To find the scale factors, we first find expressions for dx, dy and dz :

$$dx = v \cos \theta \, du + u \cos \theta \, dv - u \, v \sin \theta \, d\theta$$
$$dy = v \sin \theta \, du + u \sin \theta \, dv + u \, v \cos \theta \, d\theta$$
$$dz = u \, du - v \, dv$$

Well, we can write out lots of terms and add, or we can be appropriately lazy :

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\begin{aligned} &\ln[249]:= \mbox{Clear}[dx, dy, dz, du, dv, d\theta, u, v, \theta] \\ &dx = v \cos[\theta] du + u \cos[\theta] dv - u v \sin[\theta] d\theta; \\ &dy = v \sin[\theta] du + u \sin[\theta] dv + u v \cos[\theta] d\theta; \\ &dz = u du - v dv; \\ &Expand[dx^2 + dy^2 + dz^2] \end{aligned}
\begin{aligned} &Out[253]= du^2 u^2 - 2 du dv u v + dv^2 v^2 + dv^2 u^2 \cos[\theta]^2 + 2 du dv u v \cos[\theta]^2 + du^2 v^2 \cos[\theta]^2 + d\theta^2 u^2 v^2 \cos[\theta]^2 + d\theta^2 u^2 \sin[\theta]^2 + 2 du dv u v \sin[\theta]^2 + du^2 v^2 \sin[\theta]^2 + d\theta^2 u^2 v^2 \sin[\theta^2 + d\theta^2 u^2 \sin[\theta^2 + d\theta^2 u^2 \sin[\theta^2 + d\theta^2 + d\theta^2 u^2 \sin[\theta^2 + d\theta^2 + d\theta^2 u^2 \sin[\theta^2 + d\theta^2 +
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That's ugly; let's simplify :

In[254]:= Simplify[%]

 $\text{Out} [254]= \ d\Theta^2 \ u^2 \ v^2 \ + \ du^2 \ \left( u^2 \ + \ v^2 \right) \ + \ dv^2 \ \left( u^2 \ + \ v^2 \right)$ 

Ok, no cross terms, it's an orthogonal transformation, and the scale factors are :

$$h_u = h_v = \sqrt{u^2 + v^2} \qquad h_\phi = u v$$

To find unit vectors, write the position vector

$$\mathbf{r} = \mathbf{x}\,\hat{\mathbf{x}} + \mathbf{y}\,\hat{\mathbf{y}} + \mathbf{z}\,\hat{\mathbf{z}} = \mathbf{u}\,\mathbf{v}\cos\theta\,\hat{\mathbf{x}} + \mathbf{u}\,\mathbf{v}\sin\theta\,\hat{\mathbf{y}} + \frac{1}{2}\left(\mathbf{u}^2 - \mathbf{v}^2\right)\hat{\mathbf{z}}$$

and each unit vector is determined from :

$$\hat{\mathbf{q}}_{\mathbf{i}} = \frac{\partial \mathbf{r}}{\partial q_{\mathbf{i}}} / \left| \frac{\partial \mathbf{r}}{\partial q_{\mathbf{i}}} \right|$$

where  $\hat{q}_i$  is the generalized unit vector. Applying this definition:

$$\hat{\mathbf{u}} = \frac{\mathbf{v}\cos\theta\,\hat{\mathbf{x}} + \mathbf{v}\sin\theta\,\hat{\mathbf{y}} + \mathbf{u}\,\hat{\mathbf{z}}}{\sqrt{\mathbf{u}^2 + \mathbf{v}^2}}$$
$$\hat{\mathbf{v}} = \frac{\mathbf{u}\cos\theta\,\hat{\mathbf{x}} + \mathbf{u}\sin\theta\,\hat{\mathbf{y}} - \mathbf{v}\,\hat{\mathbf{z}}}{\sqrt{\mathbf{u}^2 + \mathbf{v}^2}}$$
$$\hat{\boldsymbol{\theta}} = \frac{-\mathbf{u}\,\mathbf{v}\sin\theta\,\hat{\mathbf{x}} + \mathbf{u}\,\mathbf{v}\cos\theta\,\hat{\mathbf{y}}}{\mathbf{u}\,\mathbf{v}} = -\sin\,\hat{\mathbf{x}} + \cos\theta\,\hat{\mathbf{y}}$$

And it is easy to take dot products of the vectors to show that the basis vectors are orthogonal.

3. See solution in separate link.

4. We make the substitutions :

$$\nu = \frac{c}{\lambda}$$
 and  $d \vee = \frac{-c}{\lambda^2} d\lambda$ 

to obtain :

$$B_{\lambda}(T) d\lambda = 2h \frac{\left(\frac{c}{\lambda}\right)^{3}}{c^{2}} \frac{1}{e^{h c/\lambda k T} - 1} \frac{c}{\lambda^{2}} d\lambda = 2h \frac{c^{2}}{\lambda^{5}} \left(e^{h c/\lambda k T} - 1\right)^{-1} d\lambda$$

Using Mathematica :

 $In[255]:= Clear[b, \lambda, t, h, c, k]$   $b[\lambda_{-}, t_{-}] := b[\lambda, t] = (2hc^{2}/\lambda^{5}) (Exp[hc/(\lambda kt)] - 1)^{-1}$   $D[b[\lambda, t], \lambda]$   $Out[257]= \frac{2c^{3}e^{\frac{ch}{kt\lambda}}h^{2}}{\left(-1 + e^{\frac{ch}{kt\lambda}}\right)^{2}kt\lambda^{7}} - \frac{10c^{2}h}{\left(-1 + e^{\frac{ch}{kt\lambda}}\right)\lambda^{6}}$ 

If you set this equal to zero and cancel common factors, you obtain :

$$\frac{\frac{h c}{\lambda k t} e^{h c/\lambda k t}}{e^{h c/\lambda k t} - 1} = 5$$

If we set  $x = h c/\lambda k t$ , this equation becomes :

$$\frac{x e^x}{e^x - 1} = 5$$

Which can be solved either via :

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In[259]:= Solve[x Exp[x] == 5 (Exp[x] - 1), x] // N
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Solve::ifun :

Inverse functions are being used by Solve, so some solutions may not be found; use Reduce for complete solution information.  $\gg$ 

 $\texttt{Out[259]=} \hspace{0.1 in} \{\hspace{0.1 in} x \hspace{0.1 in} \rightarrow \hspace{0.1 in} 0.\hspace{0.1 in} \} \hspace{0.1 in} , \hspace{0.1 in} \{\hspace{0.1 in} x \hspace{0.1 in} \rightarrow \hspace{0.1 in} 4.\hspace{0.1 in} 96511 \hspace{0.1 in} \} \hspace{0.1 in} \}$ 

or via :

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ln[261] = FindRoot[x Exp[x] = 5 (Exp[x] - 1), \{x, 5\}]
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 $\texttt{Out[261]= } \{ \texttt{x} \rightarrow \texttt{4.96511} \}$ 

Using standard MKS values for h, c and k and this value of x allows us to write Wien's Law :

$$x = \frac{hc}{\lambda k T} \Rightarrow \lambda_{max} = \frac{hc}{x k T} = \frac{0.0029}{T}$$

where T is measured in Kelvins and  $\lambda$  in meters. The sun's radiating temperature is approximately 5800 K, so the wavelength of maximum energy emission is 500 nm (5  $10^{-7}$ m)